# Defect Profiling with Low Energy Positrons of Nitrogen Implanted Silicon

D.P. van der Werf<sup>1</sup>, A.S. Saleh<sup>1</sup>, A. Towner<sup>1</sup>, M. Nathwani<sup>1</sup>, J. Taylor<sup>1</sup>, P.C. Rice-Evans<sup>1</sup> and S.J. Bull<sup>2</sup>

<sup>1</sup> Physics Department, Royal Holloway, University of London Egham, Surrey TW20 0EX, United Kingdom

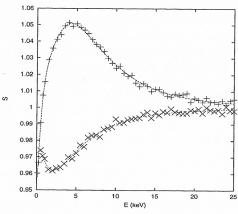
<sup>2</sup> Department of Mechanical, Materials and Manufacturing Engineering University of Newcaste Newcastle Upon Tyne, NE1 7R, United Kingdom

Keywords: Silicon, Si Defect, Nitrogen Implanted

Abstract The vacancy profile in Czochralski silicon (111) implanted with 50 keV nitrogen ions been determined using positron annihilation spectroscopy. The nitrogen distribution been measured using secondary ion mass spectroscopy. The fitted defect distribution compares well with the results of TRIM calculations.

## Introduction

Ion implantation beneath solid surfaces is widely used for sample preparation and modification science and technology. The understanding of implanted nitrogen donor impurity is especially important since the nitrogen doping of float zone silicon is used to prevent thermal slip and warpage of water subjected to high temperature processing [1,2]. Nitrogen implantation of silicon usually leads to a lager with a very high concentration of defects beneath the surface. Various techniques have been used



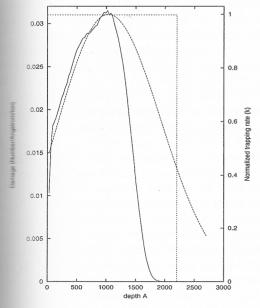
ROYDEPPROF program.

investigate the damage created during implantation and the subsequent annealing behaviour in silicon. these, positron annihilation spectroscopy (PAS) been applied successfully to study implantation will hydrogen, boron, arsenic, phosphor, oxygen, helium and fluorine [3-8]. This study extends the use of positrons to study implantation with nitrogen.

## Experimental

Czochralski (111) Si wafers were implanted with 50 keV nitrogen ions at room temperature with Kuffman ion beam. The beam consisted of 10% N and 90% N<sup>+</sup>. At the surface the N<sub>2</sub><sup>+</sup> ions immediately

Figure 1. The Doppler parameter S plotted as break into N+ and No which share the incident energy a function of positron impact energy for an First, one sample was implanted with a dose of unimplanted silicon wafer (crosses) and for 3.5×10<sup>16</sup> ions/cm<sup>2</sup>. Subsequently, the profiles of silicon implanted with 3.5×10<sup>16</sup> ions/cm<sup>2</sup> at 50 defects beneath the surface of the samples were keV (plusses). The lines are the result of fitting studied with a variable low energy positron beam the diffusion model to each set of data using (TACITUS). Then, the nitrogen profile of the sample determined using Secondary Ion



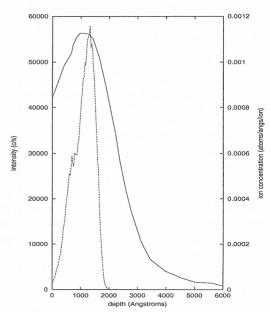


Figure 2. The defect distribution as calculated using TRIM (solid line), the fitted truncated Gaussian trapping rate (broken line) and the fixed square defect distribution.

Figure 3. The SIMS depth profile of ions for the  $3.5 \times 10^{16}$  ions/cm<sup>2</sup> implanted silicon sample (solid line) and the profile as calculated using TRIM (broken line).

spectroscopy (SIMS) [9]. The vacancy and nitrogen distributions for 50 keV nitrogen ions in silicon was calculated using the Monte Carlo ion implantation program TRIM-95 [10]. Afterwards three silicon amples were implanted with doses of 10<sup>14</sup>, 10<sup>15</sup>, 10<sup>16</sup> ions/cm<sup>2</sup>, respectively, and studied using TACITUS.

## Bults and discussion

1 shows the measured S parameter as a function of incident positron energy for silicon implanted  $3.5 \times 10^{16}$  ions/cm<sup>2</sup> and for an unimplanted specimen of the same wafer. In silicon, at room perature single vacancies are believed to disappear resulting in divacancies and for these a specific sitron trapping rate σ of  $2 \times 10^{-8}$  cm<sup>3</sup>/sec has been reported [11]. We fitted defect distributions to the curve using our own positron diffusion model program ROYDEPPROF which may accommodate any stribution of defects. The analysis of the unimplanted sample resulted in a diffusion length of 2200 Å ich is in accordance with the value of 2150 Å reported by Schultz et al. [12]. Using a square defect stribution we find a defect region of 2200 Å depth and a defect concentration  $C_v$ = 8×10<sup>-5</sup> defects per m. Using a combination of two truncated Gaussians the maximum defect concentration yields the value. One observes for TRIM that the majority of the vacancies produced lies in a region of 2000 wide with a maximum at 1000 Å. This is in poor agreement with either the square distribution or the material Gaussian. Previously, comparisons between TRIM and positron annihilation spectroscopy [13-5] showed that TRIM underestimates the range of the defect to a greater degree, except for phosphorus melanted into silicon in which compensation may be achieved by adding an electric field [16].

The nitrogen contents for the same sample were measured using Secondary Ion Mass sectroscopy. The results are shown in figure 3, together with TRIM results for the for 50 keV nitrogen implanted into silicon. Figure 3 shows that the measured peak of the nitrogen profile is at about Å and it falls to zero at about 4000 Å. The figure also shows the nitrogen ion range estimated with

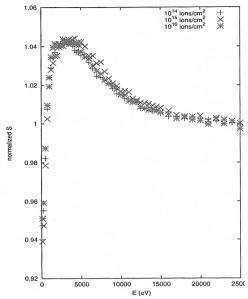


Figure 4. The Doppler parameter S plotted as a function of positron impact energy for three Acknowledgments silicon implanted with 50 keV  $N_2$  ions with different dose.

TRIM at 50 keV. It turns out that the peak position calculated by TRIM reproduces the measured value fairly well, but the width of the peak is smaller.

In figure 4 we have plotted the measured a parameter as a function of incident positron energy silicon implanted with 1014, 1015 and 1016 ions/cm is seen that the curves are almost indistinguishable low dose 1014 appears to create the same damage as a 10<sup>16</sup> exposure. In principle the value of S is linear related to the defect concentration C<sub>v</sub>. Thus a dose might be expected to cause greater damage (higher C<sub>v</sub>) which would be indicated by a higher 5 This is not the case here. The equal heights in S at 2 keV might suggest that all the positrons are bear trapped in both cases; i.e we have a saturation positron trapping. However, the left hand side of curve yields a diffusion length of 500 Å which is large than expected for positron saturation. Hence, we conclude that we may have defect saturation, i.e. defect concentration is independent of the dose.

The research was supported by the Engineering Physical Science Research Council (grant number

GR/K 30964) and the Corporate Research Programme of AEA Technology.

### References

- [1] R.J. Dexter, S.B. Wateski and S.T. Picraux, App. Phys. Lett 32 (1973) 455
- [2] T. Tsujide, M. Nojiri and H. Kitagawa, J. App. Phys 51 (1980) 1605
- [3] J. Keinonen, M. Hautala, E. Rauhala, V. Karttunen, A. Kuronen, J. Räisänen, J. Lattinen. Lahtinen, A. Vehanen, E. Punaak and P. Hautojäarvi, Phys. Rev. B 37 (1988) 8269
- [4] D. Smith, H. Evans, C. Smith, P.C. Rice-Evans and J. Evans in: Kajcsos Ks and Szeles (eds), Proc 9th int. Conf. on Positron Annihilation, Tran Tech Publications, Aedermannsdorf (1992) 1451
- [5] A. Uedono, S. Tanigawa, J. Sugiura and M. Ogasawara, Jpn. J. Appl. Phys. 29 (1990) 1867
- [6] P. Simpson, M. Vos, I. Mitchell, C. Wu and P. Schultz, Phys. Rev. B 44 (1991) 12180
- [7] P. Simpson, M. Vos, I. Mitchell, C. Wu and P. Schultz in: Kajcsos Ks and Szeles (eds), Proc. 9th interpretations of the control of the co Conf. on Positron Annihilation, Tran Tech Publications, Aedermannsdorf (1992) 1439
- [8] M. Fujinami and N. Chilton, J. Appl. Phys. 73 (1993) 3242
- [9] A. Benninghoven, F. Rudenauer and H. Werner, Secondary Ion mass Spectrometry, Basic Concepts Instrumental Aspects, applications and Trends. Wiley, New York (1987).
- [10] J. Ziegler, J. Biersack and U. Litimark, The Stopping and Range of Ions in Solids, Vol. (PergamonPress, New York, N., Y. 1985)
- [11] P. Mascher, S. Dannefaer and D. Kerr, Phys. Rev. B40 (1989) 11764
- [12] P.J. Schultz, E. Tandberg, B. Nielsen, T.E. Jackman and K.G. Lynn, Phys. Rev. Lett. 61 (1988) 187
- [13] D. Smith, P.C. Rice-Evans, D. Britton, J. Evans and A. Allen, Phil. Mag.A 61 (1990) 839
- [14] T. Aruga, S. Takamura, M. Hirose and Y. Itoh, Phys. Rev. B46 (1992) 14411
- [15] A. Uedono, S. Tanigawa and H. Sakairi, J. Nucl. Mater. 184 (1991) 191
- [16] P. Asoka-Kumar, P. Sferlazzo, H. Au and K. Lynn, Nucl. Inst. Metho. B74 (1993) 89